COSMIC MAGNETIC FIELDS

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fluid dynamics

CONSERVATION OF PARTICLE NUMBER AND MASS

- Consider a (Lagrangian) volume element δV containing δN particles. The particle number density is $n = \frac{\delta N}{\delta V}$.
- Unless particles can be created or destroyed, the particle number is conserved d(δN) = 0.
- For particles of mass *m* the mass element $\delta M = m \,\delta N$ is conserved $d(\delta M) = 0$. Mass density is $\rho = \frac{\delta M}{\delta V} = mn$.
- A volume element evolves according to the local velocity field $\vec{v}(\vec{r})$. During a time interval d*t* it will be displaced by $d\vec{r} = \vec{v} dt$ and expand by $d(\delta V) = \delta V(\vec{\nabla} \cdot \vec{v}) dt$.

• With the material derivative:
$$\frac{\partial n}{\partial t} + \left(\vec{v} \cdot \vec{\nabla}\right)n = \frac{\mathrm{d}n}{\mathrm{d}t} = \frac{\mathrm{d}(\delta N/\delta V)}{\mathrm{d}t} = -n\left(\vec{\nabla} \cdot \vec{v}\right)$$

Continuity equation: $\frac{\partial n}{\partial t} + \vec{\nabla} \cdot (n\vec{\nabla}) = 0, \quad \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho\vec{\nabla}) = 0$

CONSERVATION OF MOMENTUM

• Consider a force $\delta \vec{F}$ acting on a fluid element of volume δV , mass δM , velocity \vec{v} , and momentum $\delta \vec{p} = \vec{v} \, \delta M = \rho \vec{v} \, \delta V$.

Newton's second law:
$$\delta \vec{F} = \frac{d(\delta \vec{p})}{dt} = \frac{d \vec{v}}{dt} \delta M$$

- Euler equation: $\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla})\vec{v} = \frac{d\vec{v}}{dt} = \frac{\delta \vec{F}}{\delta M} = \frac{\vec{f}}{\rho},$ introducing the force density $\vec{f} = \frac{\delta \vec{F}}{\delta V}$
- Combined with the continuity equation, this can be written in the tensor form: $\partial_t(\rho v^i) + \partial_j(\rho v^i v^j) = f^i$

KINETIC ENERGY DENSITY

•
$$u_{\rm kin} = \frac{\delta \mathscr{E}_{\rm kin}}{\delta V} = \rho v^2 / 2$$

• Continuity and Euler equations lead to: $\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$ $\vec{v} \cdot \left[\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right] = \vec{v} \cdot \frac{\vec{f}}{\rho}$ $\frac{\partial u_{\text{kin}}}{\partial t} + \vec{\nabla} \cdot (u_{\text{kin}} \vec{v}) = \frac{du_{\text{kin}}}{dt} + u_{\text{kin}} (\vec{\nabla} \cdot \vec{v}) = \vec{f} \cdot \vec{v}$

The last term represents work done on the fluid element in unit time.

INTERNAL ENERGY DENSITY

•
$$u_{\text{int}} = \frac{\delta \mathscr{C}_{\text{int}}}{\delta V}$$

- First law of thermodynamics: $d(\delta \mathscr{C}_{int}) = T d(\delta S) P d(\delta V)$ with temperature *T*, entropy δS , pressure *P*.
- In an adiabatic process, entropy is conserved, $d(\delta S) = 0$, hence $d(\delta \mathscr{E}_{int}) = -P d(\delta V)$.
- $\frac{\mathrm{d}u_{\mathrm{int}}}{\mathrm{d}t} = -(u_{\mathrm{int}} + P)\left(\overrightarrow{\nabla}\cdot\overrightarrow{\mathbf{v}}\right)$
- $w = u_{int} + P = \kappa u_{int}$ is the enthalpy density, with κ the adiabatic index (specific heats ratio)
- For ideal non-relativistic monoatomic gas: $\kappa = 5/3$, $u_{int} = (3/2)P$ and w = (5/2)P.

•
$$\frac{\partial P}{\partial t} + \left(\overrightarrow{\mathbf{v}} \cdot \overrightarrow{\nabla} \right) P = \frac{\mathrm{d}P}{\mathrm{d}t} = -\kappa P \left(\overrightarrow{\nabla} \cdot \overrightarrow{\mathbf{v}} \right)$$

CONSERVATION OF ENERGY

- Pressure gradient as a force density: $\vec{f}_{\rm P} = \vec{\nabla} P$.
- Work done by pressure gradient changes the kinetic energy: $d(\delta \mathscr{C}_{kin}) = d(\delta W_{P}) = \delta V \left(\vec{f}_{P} \cdot \vec{v} \right) dt = -\delta V dP.$
- Internal energy change: $d(\delta \mathscr{C}_{int}) = -P d(\delta V)$.
- Total energy change: $d(\delta \mathscr{E}_{tot}) \equiv d(\delta \mathscr{E}_{kin} + \delta \mathscr{E}_{int}) = -d(P \,\delta V).$
- Conservation of total energy: $d(\delta \mathscr{C}_{tot} + P \,\delta V) = 0$.

• Bernoulli's equation:
$$\frac{d(\delta \mathscr{E}_{tot} + P \,\delta V)}{\delta M} = d\left(\frac{v^2}{2} + \frac{w}{\rho}\right) = 0.$$

equivalent tensor form: $\frac{\partial u_{\text{tot}}}{\partial t} + \vec{\nabla} \cdot \left[\left(u_{\text{kin}} + w \right) \vec{\nabla} \right] = 0.$

CONSERVATION OF ENTROPY

• Relation between density and pressure variations: $\vec{\nabla} \cdot \vec{v} = -\frac{1}{\rho} \frac{d\rho}{dt} = -\frac{1}{\kappa P} \frac{dP}{dt}$

- Adiabatic invariant: $d[\ln(P/\rho^{\kappa})] = 0$.
- Adiabatic equation of state: $P \propto \rho^{\kappa}$ for a given Lagrangian fluid element.
- Specific entropy: $s = \frac{\delta S}{\delta M} = c_V \ln(P/\rho^{\kappa})$ with $c_V = (3/2)(k_{\rm B}/m)$: the specific heat capacity at constant volume.
- ds = 0: equivalent to energy equation for adiabatic fluid.
- Conservation of mass implies the conservation of entropy $d(\delta S) = 0$.

EQUATIONS FOR FLUID DYNAMICS

• continuity equation (conservation of mass):

 $\frac{\partial \rho}{\partial t} + \overrightarrow{\nabla} \cdot \left(\rho \overrightarrow{\mathbf{v}} \right) = 0$

- Euler equation (conservation of momentum): $\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = \frac{\vec{f}}{\rho}$
- pressure equation (conservation of energy and entropy): $\frac{\partial P}{\partial t} + \left(\vec{v} \cdot \vec{\nabla}\right)P + \kappa P\left(\vec{\nabla} \cdot \vec{v}\right) = 0$

LINEARIZATION

- $\rho = \rho_0 + \rho_1$ with $|\rho_1| \ll |\rho_0|$, etc.
- continuity equation:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{\mathbf{v}}) = 0$$
$$\frac{\partial \rho_1}{\partial t} + \vec{\nabla} \cdot (\rho_1 \vec{\mathbf{v}}_0) + \vec{\nabla} \cdot (\rho_0 \vec{\mathbf{v}}_1) = 0$$

• Euler equation:

$$\frac{\partial \overrightarrow{\mathbf{v}}}{\partial t} + \left(\overrightarrow{\mathbf{v}} \cdot \overrightarrow{\nabla}\right) \overrightarrow{\mathbf{v}} = \frac{\overrightarrow{f}}{\rho}$$
$$\frac{\partial \overrightarrow{\mathbf{v}}_1}{\partial t} + \left(\overrightarrow{\mathbf{v}}_1 \cdot \overrightarrow{\nabla}\right) \overrightarrow{\mathbf{v}}_0 + \left(\overrightarrow{\mathbf{v}}_0 \cdot \overrightarrow{\nabla}\right) \overrightarrow{\mathbf{v}}_1 = \frac{\overrightarrow{f}_1}{\rho_0} - \frac{\overrightarrow{f}_0}{\rho_0^2}\rho_1$$

• pressure equation:

$$\frac{\partial P}{\partial t} + \left(\overrightarrow{v} \cdot \overrightarrow{\nabla}\right) P + \kappa P\left(\overrightarrow{\nabla} \cdot \overrightarrow{v}\right) = 0$$

$$\frac{\partial P_1}{\partial t} + \left(\overrightarrow{v}_1 \cdot \overrightarrow{\nabla}\right) P_0 + \left(\overrightarrow{v}_0 \cdot \overrightarrow{\nabla}\right) P_1 + \kappa P_1\left(\overrightarrow{\nabla} \cdot \overrightarrow{v}_0\right) + \kappa P_0\left(\overrightarrow{\nabla} \cdot \overrightarrow{v}_1\right) = 0$$

LINEARIZATION: UNIFORM STATIC BACKGROUND

• Assume static ($\vec{v}_0 = 0$) and uniform (constant ρ_0, P_0) background. This implies that $\vec{f}_0 = -\vec{\nabla}P_0 = 0$.

 ρ_1

• continuity equation:

 $\frac{\partial \rho_1}{\partial t} + \vec{\nabla} \cdot \left(\rho_1 \vec{\mathbf{v}}_0\right) + \vec{\nabla} \cdot \left(\rho_0 \vec{\mathbf{v}}_1\right) = 0$ $\frac{\partial \rho_1}{\partial t} = -\rho_0 \left(\vec{\nabla} \cdot \vec{\mathbf{v}}_1\right)$

• Euler equation:

$$\frac{\partial \overrightarrow{\mathbf{v}}_{1}}{\partial t} + \left(\overrightarrow{\mathbf{v}}_{1} \cdot \overrightarrow{\nabla}\right) \overrightarrow{\mathbf{v}}_{0} + \left(\overrightarrow{\mathbf{v}}_{0} \cdot \overrightarrow{\nabla}\right) \overrightarrow{\mathbf{v}}_{1} = \frac{\overrightarrow{f}_{1}}{\rho_{0}} - \frac{\overrightarrow{f}_{0}}{\rho_{0}^{2}}\rho_{1}$$
$$\frac{\partial \overrightarrow{\mathbf{v}}_{1}}{\partial t} = -\frac{\overrightarrow{\nabla}P_{1}}{\rho_{0}}$$

• pressure equation:

$$\frac{\partial P_1}{\partial t} + \left(\vec{\mathbf{v}}_1 \cdot \vec{\nabla}\right) P_0 + \left(\vec{\mathbf{v}}_0 \cdot \vec{\nabla}\right) P_1 + \kappa P_1 \left(\vec{\nabla} \cdot \vec{\mathbf{v}}_0\right) + \kappa P_0 \left(\vec{\nabla} \cdot \vec{\mathbf{v}}_1\right) = 0$$

$$\frac{\partial P_1}{\partial t} = -\kappa P_0 \left(\vec{\nabla} \cdot \vec{\mathbf{v}}_1\right)$$

LINEARIZATION: UNIFORM STATIC BACKGROUND

• Adopt oscillatory velocity perturbation $\vec{v}_1 \propto \exp\left(i\omega t + i\vec{k}\cdot\vec{r}\right)$

• pressure equation:

 $\frac{\partial P_1}{\partial t} = -\kappa P_0 \left(\vec{\nabla} \cdot \vec{\mathbf{v}}_1 \right)$ $i\omega P_1 = -\kappa P_0 \left(i\vec{k} \cdot \vec{\mathbf{v}}_1 \right) \quad P_1 \propto \exp\left(i\omega t + i\vec{k} \cdot \vec{r} \right)$

• Euler equation:

$$\frac{\partial \vec{v}_1}{\partial t} = -\frac{\vec{\nabla} P_1}{\rho_0}$$
$$i\omega \vec{v}_1 = -\frac{i\vec{k} P_1}{\rho_0} = \frac{i\vec{k}}{\rho_0} \frac{\kappa P_0}{\omega} \left(\vec{k} \cdot \vec{v}_1\right): \text{ longitudinal velocity perturbation } (\vec{v}_1 \parallel \vec{k}), \text{ hence } P_1 \neq 0 \text{ (compression).}$$

- dot product with $-i\omega \vec{k}$ yields: $\omega^2 \left(\vec{k} \cdot \vec{v}_1\right) = k^2 \frac{\kappa P_0}{\rho_0} \left(\vec{k} \cdot \vec{v}_1\right)$
- dispersion relation $\frac{\omega^2}{k^2} = \frac{\kappa P_0}{\rho_0} \equiv c_{s,0}^2$: stable ($\omega^2 > 0$) wave the sound wave.

 $c_{s,0}$ is both the phase and group speed (uniform and isotropic) – the speed of sound.

RELATIVISTIC SPEED OF SOUND

• Maxwell-Jüttner distribution $f(\gamma) d\gamma = \frac{\gamma^2 \beta d\gamma}{\Theta K_2(1/\Theta)} \exp\left(-\frac{\gamma}{\Theta}\right)$ where $\Theta = k_{\rm B} T/mc^2$ is the relativistic temperature, $\gamma = (1 - \beta^2)^{-1/2}$ is the particle Lorentz factor, $K_2(x)$ is the modified Bessel function of the second kind.

mean particle energy $\langle \gamma \rangle = 3\Theta + h(\Theta)$ with $h(\Theta) = \frac{K_1(1/\Theta)}{K_2(1/\Theta)}$ pressure $P = nk_{\rm B}T = \Theta\rho c^2$ relativistic enthalpy density $w = \rho c^2 + \frac{\kappa}{\kappa - 1}P = [4\Theta + h(\Theta)]\rho c^2$ adiabatic index $\kappa = 1 + \frac{\Theta}{\langle \gamma \rangle - 1}$ speed of sound $c_{\rm s} = c\sqrt{\frac{\kappa P}{w}}$



In the limit of non-relativistic temperatures $\Theta \ll 1$: $h(\Theta) \simeq 1 - \frac{3}{2}\Theta$, $\kappa \simeq \frac{5}{3}$, $w \simeq \rho c^2 + \frac{5}{2}P$,

$$c_{\rm s} \simeq c \sqrt{\frac{5\Theta}{3}} = \sqrt{\frac{5P}{3\rho}}$$

• In the limit of ultra-relativistic temperatures $\Theta \gg 1$: $h(\Theta) = \frac{1}{2\Theta}, \kappa \simeq \frac{4}{3}, w \simeq 4P, c_s \simeq \frac{c}{\sqrt{3}}$.