

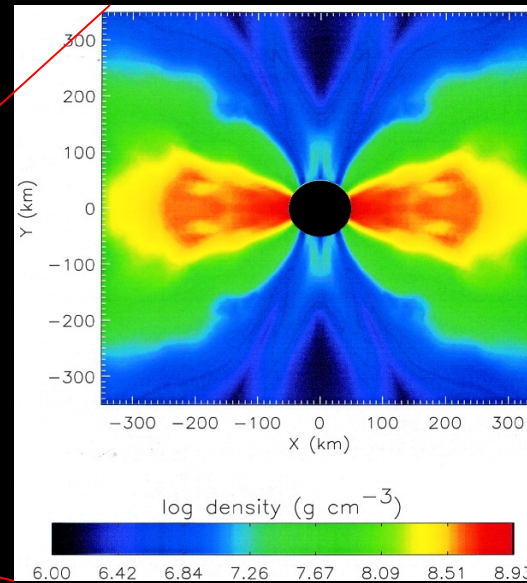
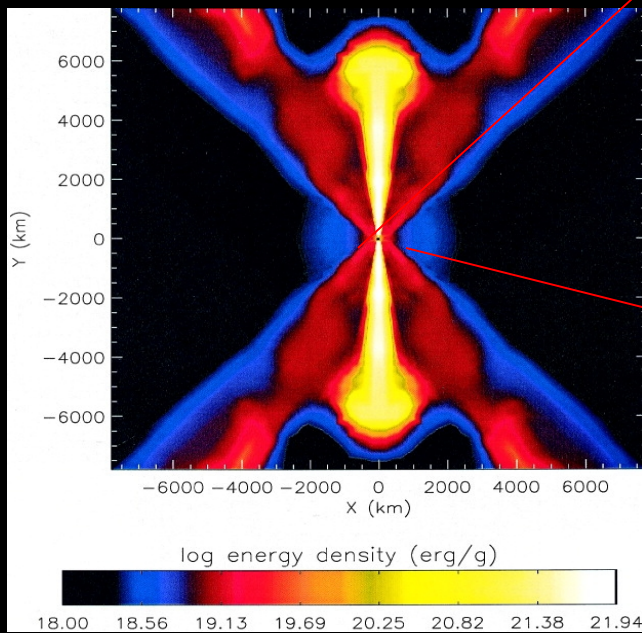
# Accretion in the collapsar. How long can be a long GRB?

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# Collapsar in the engine of long GRBs



The thick torus may form during the collapse of a massive, rotating star. This torus is likely to be connected with a jet production and gamma ray burst.

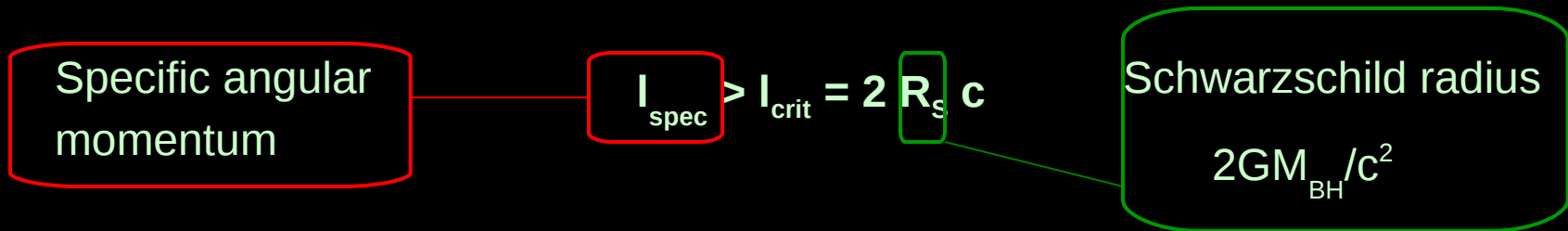
Simulations done by MacFadyen & Woosley (1999)

Mechanisms of energy transfer from torus to jet:

1. neutrinos
2. BH rotation

# Conditions for torus existence

The rotation must prevent the envelope material from the radial infall onto BH.



**BUT:** BH is growing fast (accretion rate of 0.01-1 Msun/s) =>  
the GRB is emitted only until  $I > I_{\text{crit}}$  is satisfied.

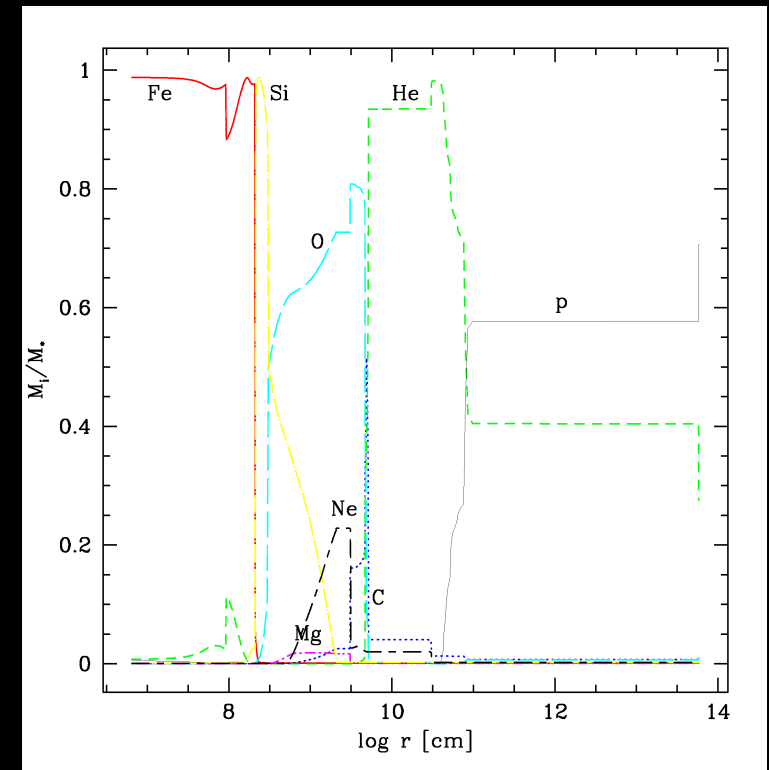
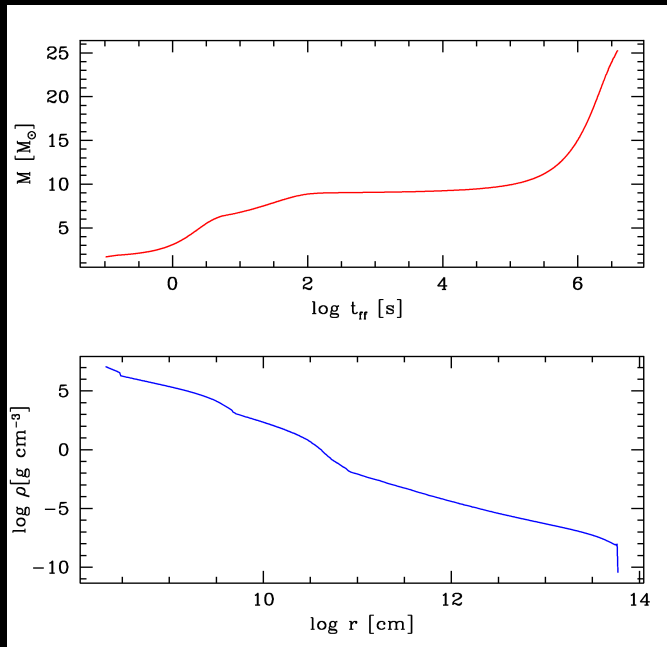
For a rotating BH:

$$I_{\text{crit}} = 2GM_{\text{BH}}/c (2-A+2(1-A)^{1/2})^{1/2}$$

*How long can be a GRB?*

# Pre-SN model: iron core + envelope

Density and mass distribution in the pre-supernova star



**Chemical composition**  
(data from Woosley & Weaver 1995)

# How the pre-collapse star rotates?

The distribution of specific angular momentum in the pre-SN star unknown.

Some assumptions:

Polar angle dependence (differential rotation)

Radius dependence (rigid rotation, with a possible cut-off on  $l_{\text{spec}}$ )

Constant ratio of centrifugal to gravitational forces

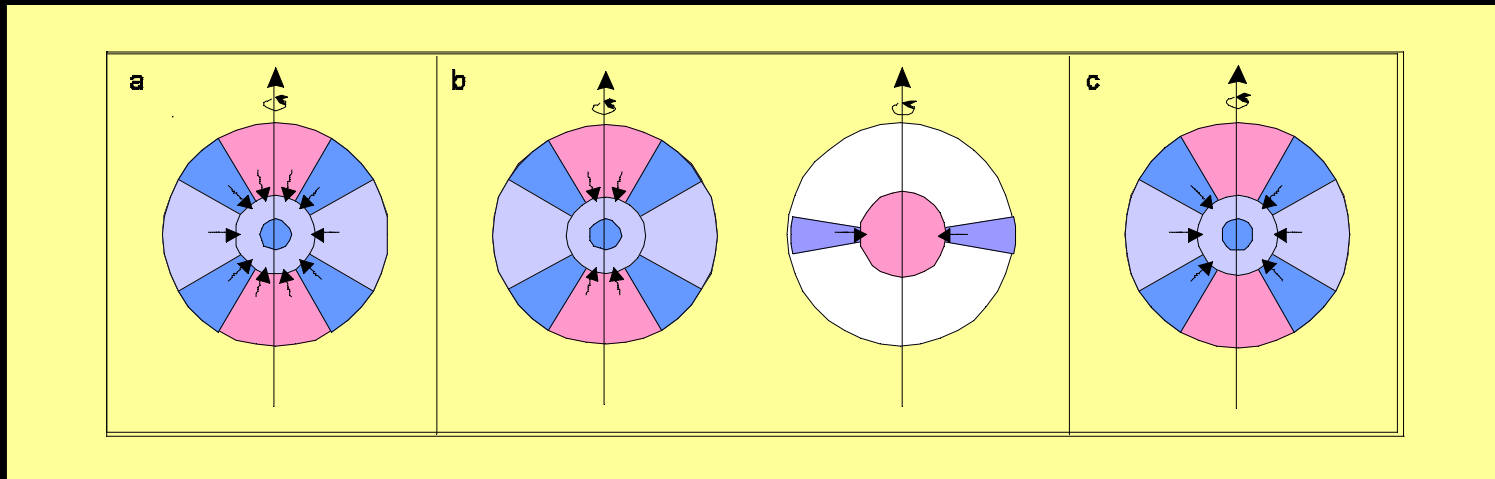
$$l_{\text{spec}} = l_0 (1 - \cos \theta)$$

$$l_{\text{spec}} = l_0 (r/r_{\text{in}}) \sin^2 \theta$$

# The black hole grows due to accretion

The time evolution of the collapsar => iterative procedure

1. BH mass = iron core mass
2. Envelope shells accrete
3. Check for conditions given by the changing BH mass and spin

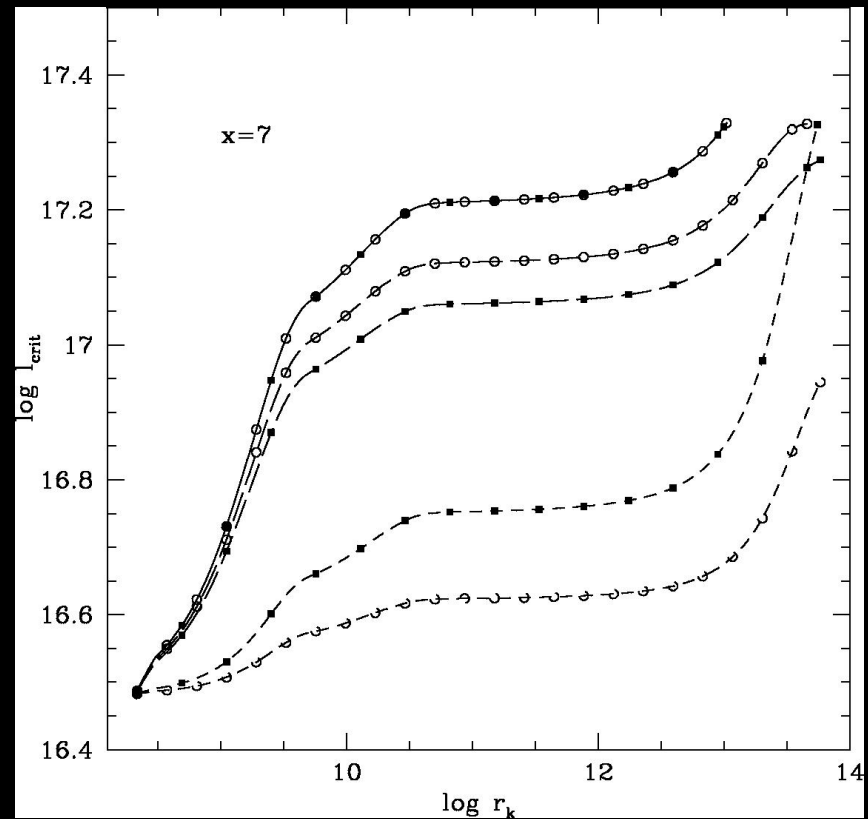


Various possible accretion scenarios

# Critical angular momentum during the collapse

Large  $l_{\text{crit}}$   $\Rightarrow$  less material in the envelope can form torus

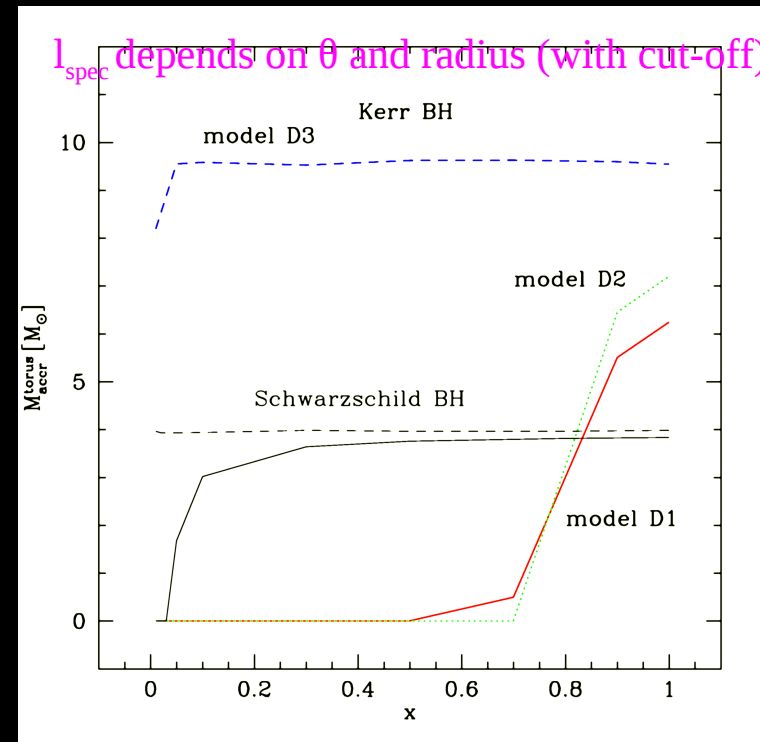
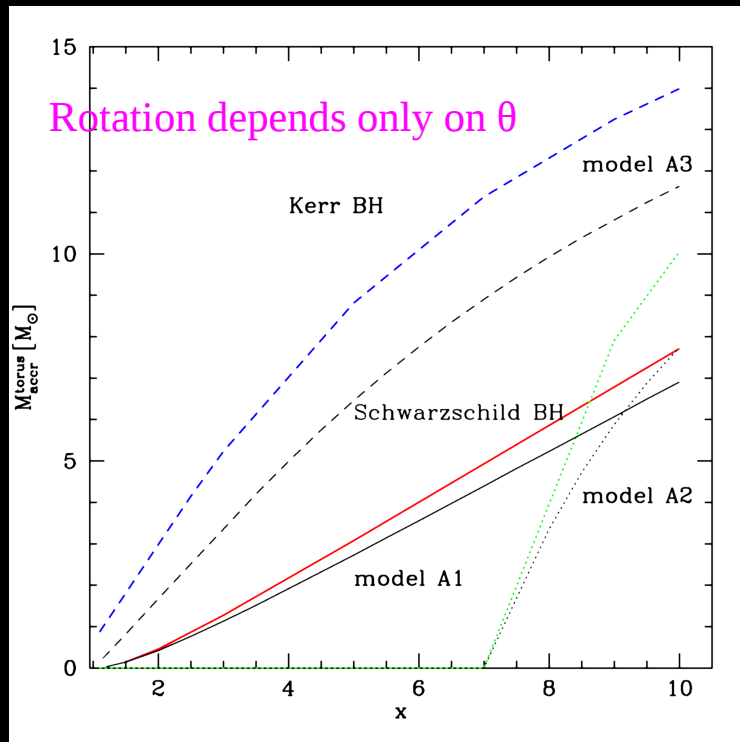
No torus  $\Rightarrow$  GRB is finished (the jet is accretion powered)



Schwarzschild BH case (Janiuk & Proga, 2008, ApJ, 675, 519)

# How long is a GRB?

Constant accretion rate => GRB duration proportional to mass accreted through torus

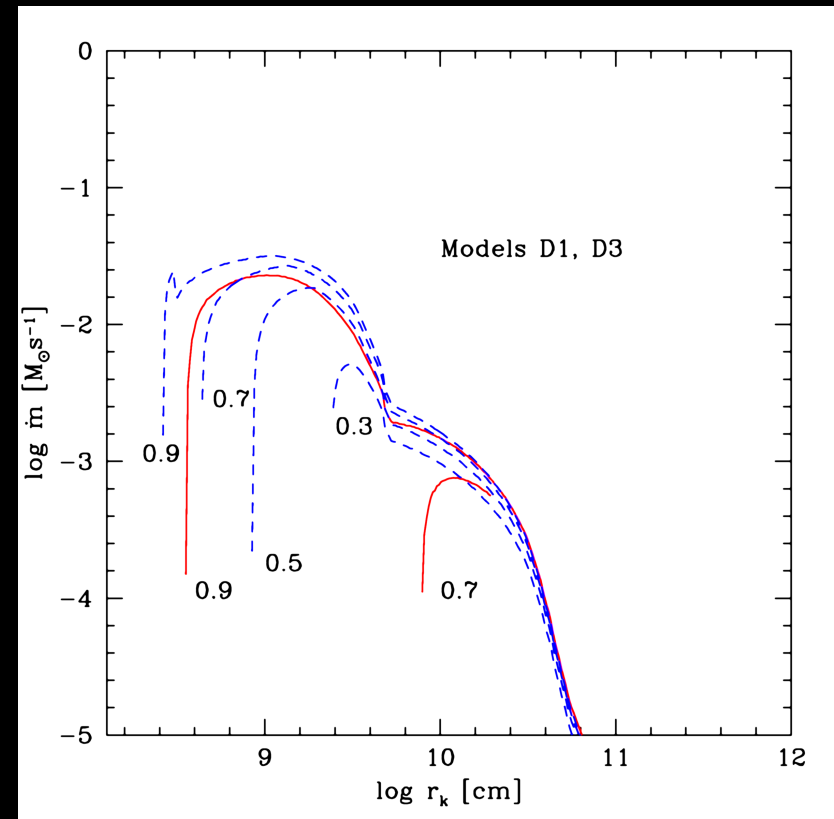
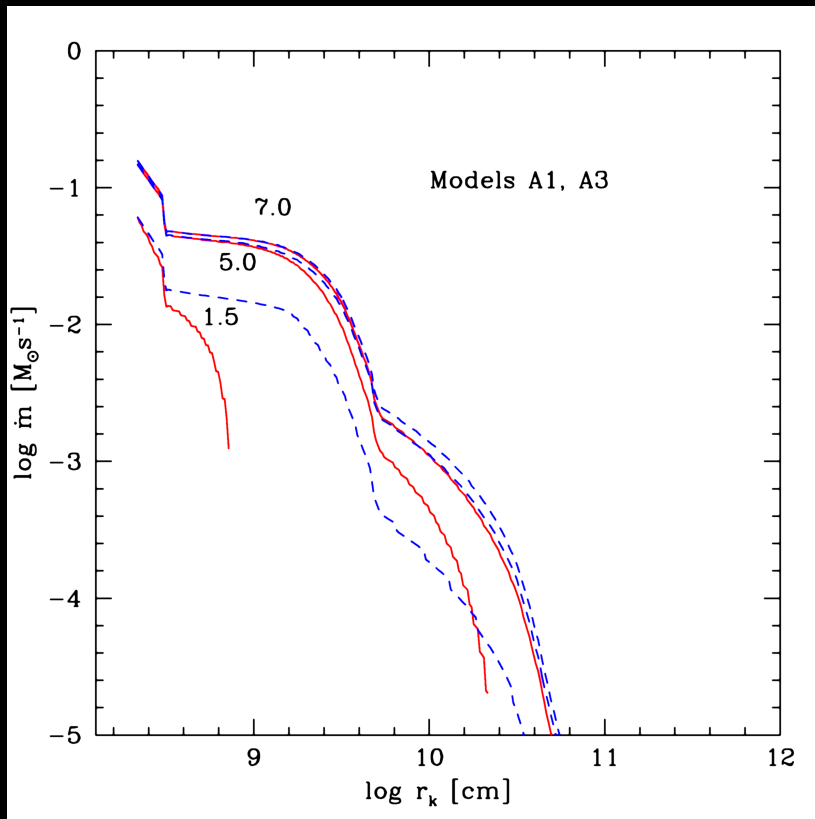


Accreted mass as a function of initial normalization of  $l_{\text{spec}}$  and for various models of rotation.

*Kerr BH case: Janiuk, Moderski & Proga, 2008, ApJ, submitted*



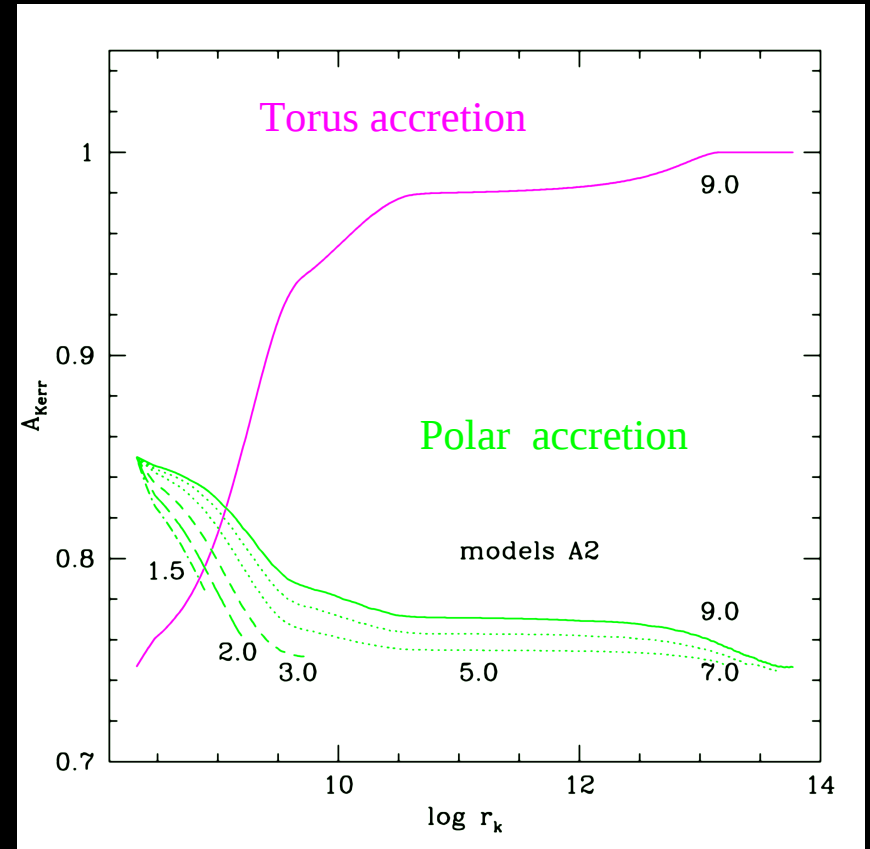
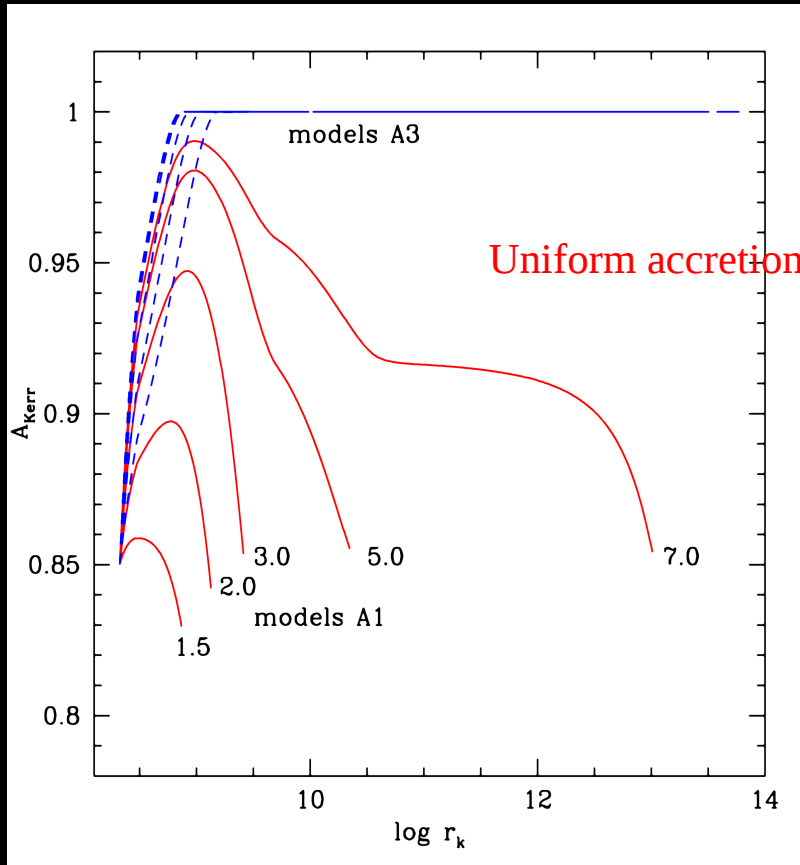
# Accretion rate is not constant



*Neutrino annihilation inefficient* when accretion rate  $< 0.01 M_{\text{sun}} / \text{s}$   $\Rightarrow$  GRB finishes

(e.g. Popham et al. 1999; Di Matteo et al. 2002; Janiuk et al. 2004; 2007)

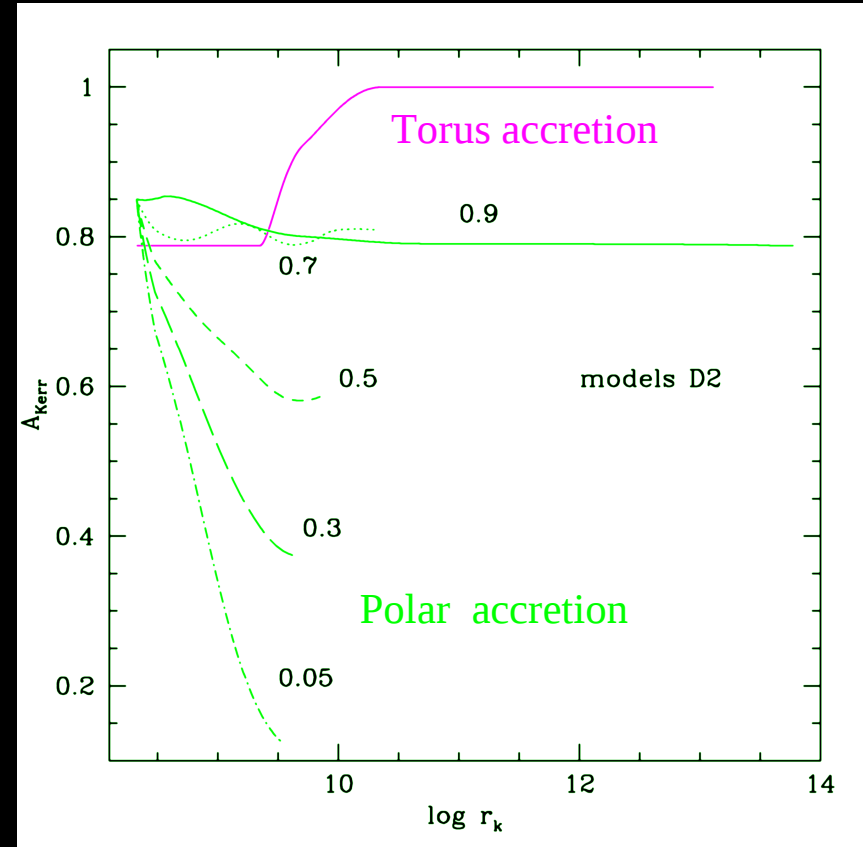
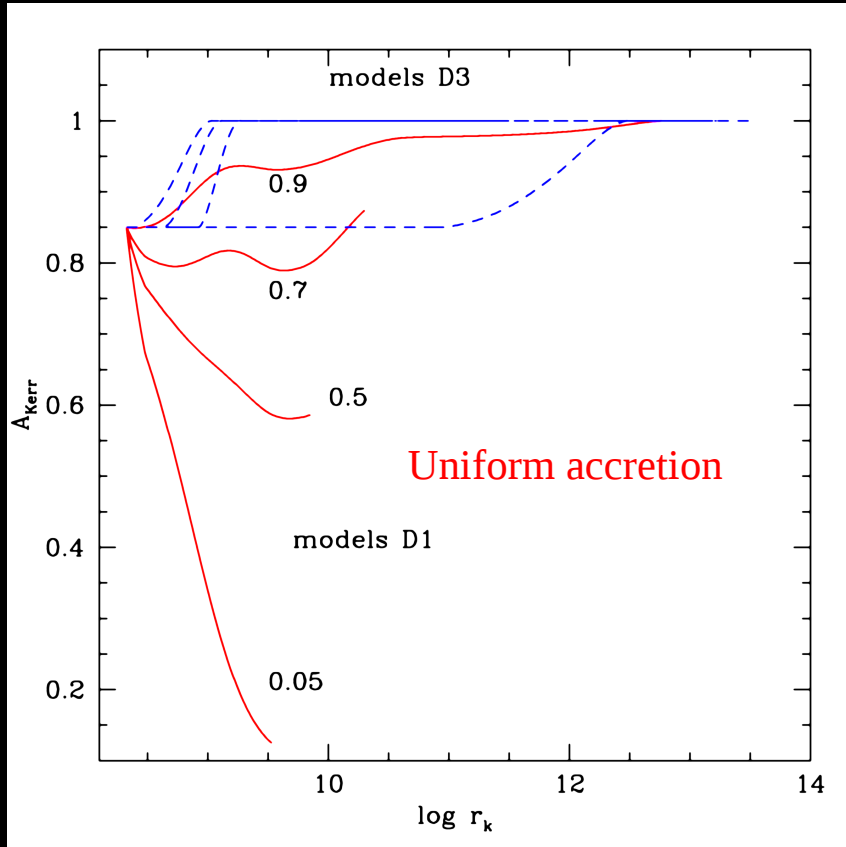
# Evolution of the BH spin



Differential rotation models

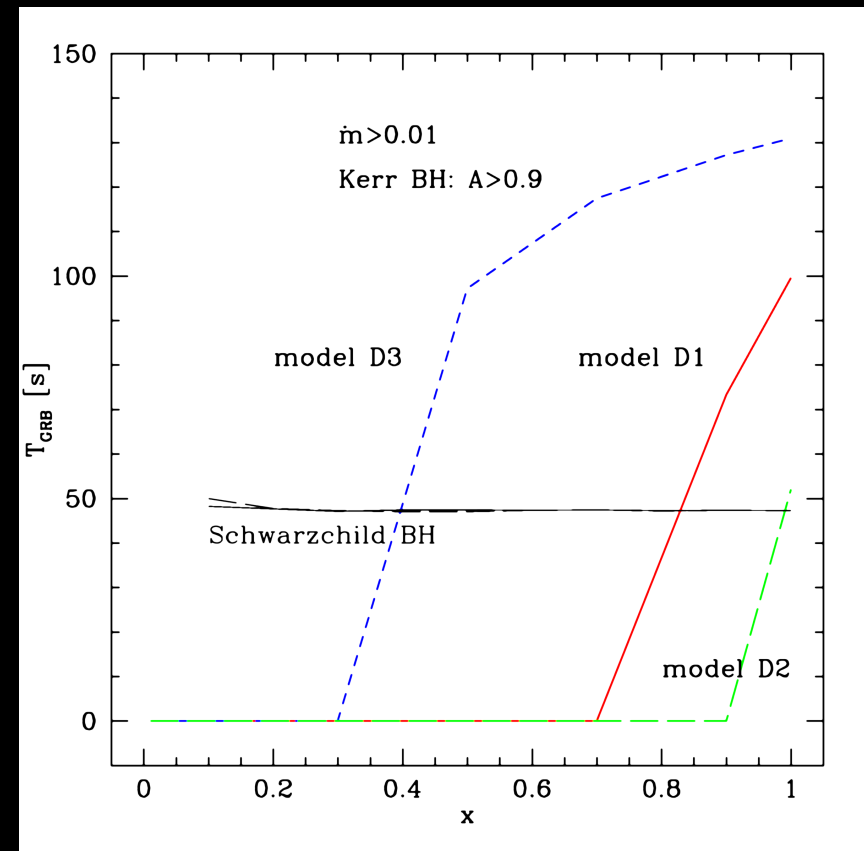
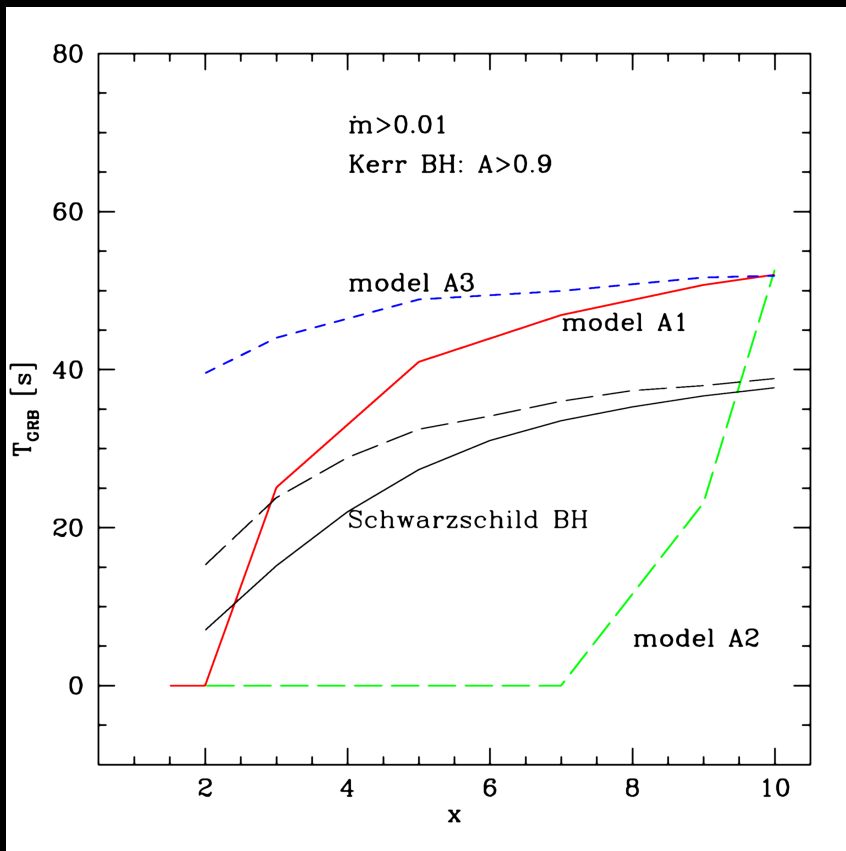
*B-Z mechanism inefficient when  $A < 0.9$  (Blandfor-Znajek 1977; McKinney 2005)*

# ... BH may even spin down completely

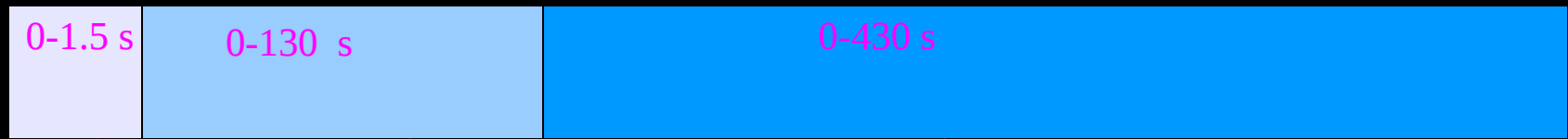


Rotation models with centrifugal-gravity balance

# GRB: large accretion rate *and* spinning BH



# Three kinds of jets in the collapsar



*Precursor jet,  
powered by  $\nu\bar{\nu}$   
large  $\dot{m}$ , small  $A$*

*First jet, powered  
by both  $\nu\bar{\nu}$  and  
B-Z  
large  $\dot{m}$ , large  $A$*

*Second jet, powered  
by B-Z  
small  $\dot{m}$ , large  $A$*

# Summary

GRB long durations may provide constraints for the rotation law in the pre-SN star.

Realistic rotation laws result in mass accreted through the torus less than 4-15  $M_{\text{sun}}$ .

The minimum accretion rate limit for the *neutrino-powered jets*, in the Schwarzschild black hole models, results in GRB durations up to 40-100 s.

The minimum *accretion rate and BH spin* limit, for jets powered by both neutrinos and B-Z mechanism, results in GRB durations up to 50-130 s.

In the Kerr black hole models, we find the solutions corresponding to three kinds of jets: precursor jet, early jet and late jet, powered by different mechanisms. Possibly, the opening angle of these jets is changing, which would have strong observational consequences.